

The Schrödinger equation in 3D

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi + V \Psi(x, y, z, t)$$

$$\rightarrow E \Psi(x, y, z) = \left[-\frac{\hbar^2}{2m} \nabla^2 + V(x, y, z) \right] \Psi(x, y, z)$$

(time-indep. form)

$$\vec{p} = \frac{\hbar}{i} \nabla \quad \text{momentum operator}$$

$$\nabla \equiv \left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$$

$$\nabla^2 \equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

i) Free particle ($V=0$)

$$\Psi_{\vec{k}}(x, y, z) = e^{i\vec{k} \cdot \vec{r}}$$

$$\vec{r} = (x, y, z), \quad \vec{k} = (k_x, k_y, k_z)$$

$$\hat{H} \psi_{\vec{k}}(\vec{r}) = \frac{\hbar^2}{2m} (k_x^2 + k_y^2 + k_z^2) \psi_{\vec{k}}(\vec{r})$$

$$= \frac{\hbar^2 \vec{k}^2}{2m} \psi_{\vec{k}}(\vec{r})$$

$$\Rightarrow E_{\vec{k}} = \frac{\hbar^2 \vec{k}^2}{2m}$$

ii) Particle in a box

$$0 < x < L_1$$

$$0 < y < L_2$$

$$0 < z < L_3$$

Separation of variables:

$$\text{let } \psi(x, y, z) = \psi_1(x) \psi_2(y) \psi_3(z)$$

$$-\frac{2mE}{\hbar^2} = \frac{1}{\psi_1} \frac{d^2\psi_1}{dx^2} + \frac{1}{\psi_2} \frac{d^2\psi_2}{dy^2} + \frac{1}{\psi_3} \frac{d^2\psi_3}{dz^2}$$

$$\frac{1}{\psi_1} \frac{d^2\psi_1}{dx^2} = \text{const.} \equiv -\frac{2mE_1}{\hbar^2}$$

$$\text{B.C.s: } \psi_1(0) = 0 = \psi_1(L_1)$$

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$$\Rightarrow \psi_1(x) = A \sin\left(\frac{n_1 \pi x}{L_1}\right),$$

$$E_1 = \frac{\hbar^2 \pi^2 n_1^2}{2m L_1^2}, \quad n_1 = 1, 2, 3, \dots, \infty$$

$$\text{similarly, } \psi_2(y) = B \sin\left(\frac{n_2 \pi y}{L_2}\right),$$

$$E_2 = \frac{\hbar^2 \pi^2 n_2^2}{2m L_2^2}, \quad n_2 = 1, 2, 3, \dots, \infty;$$

$$\psi_3(z) = C \sin\left(\frac{n_3 \pi z}{L_3}\right),$$

$$E_3 = \frac{\hbar^2 \pi^2 n_3^2}{2m L_3^2}, \quad n_3 = 1, 2, 3, \dots, \infty$$

$$\Rightarrow \psi(x, y, z) = D \sin\left(\frac{n_1 \pi x}{L_1}\right) \sin\left(\frac{n_2 \pi y}{L_2}\right) \sin\left(\frac{n_3 \pi z}{L_3}\right)$$

$$E = E_1 + E_2 + E_3 = \frac{\hbar^2 \pi^2}{2m} \left(\frac{n_1^2}{L_1^2} + \frac{n_2^2}{L_2^2} + \frac{n_3^2}{L_3^2} \right)$$

$$0 < n_1, n_2, n_3 \in \mathbb{Z} \quad (D = ABC)$$

Normalization

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$$1 = \int_0^{L_1} dx \int_0^{L_2} dy \int_0^{L_3} dz |\Psi(x, y, z)|^2$$

$$= |D|^2 \int_0^{L_1} dx \sin^2\left(\frac{n_1 \pi x}{L_1}\right) \int_0^{L_2} dy \sin^2\left(\frac{n_2 \pi y}{L_2}\right)$$

$$\times \int_0^{L_3} dz \sin^2\left(\frac{n_3 \pi z}{L_3}\right)$$

$$= |D|^2 \frac{L_1}{2} \frac{L_2}{2} \frac{L_3}{2} = |D|^2 \frac{V_{\text{box}}}{8}$$

$$\Rightarrow D = \sqrt{\frac{8}{V_{\text{box}}}}$$

iii) Special case: Cube, $L_i = L$

$$E = \frac{\hbar^2 \pi^2}{2m L^2} (n_1^2 + n_2^2 + n_3^2)$$

$$= \frac{\hbar^2 \pi^2 \vec{n}^2}{2m L^2}$$

$$\frac{2mL^2}{\hbar^2\pi^2} E$$

n_1

n_2

n_3

3

1

1

1

6

2

1

1

6

1

2

1

6

1

1

2

9

2

2

1

9

2

1

2

9

1

2

2

etc.

} $d=3$

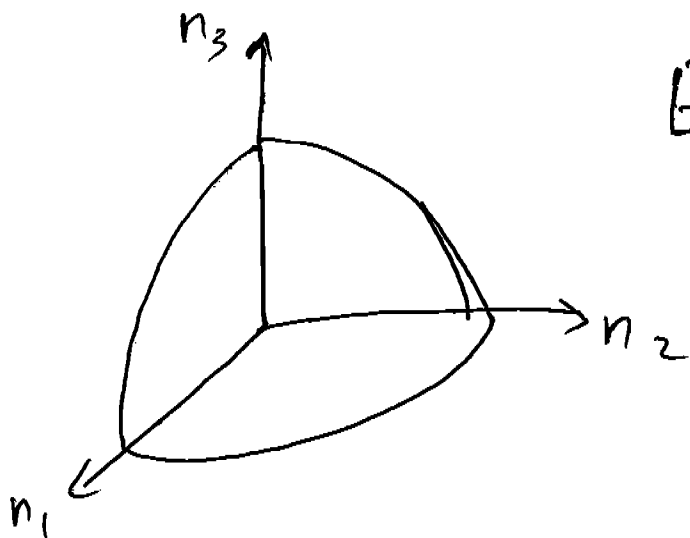
} $d=3$

Counting states

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$$N(E < E) \approx \frac{1}{8} \frac{4\pi}{3} n(E)^3 = \frac{\pi}{6} n(E)^3$$

$$E = \frac{\pi^2 \hbar^2 \vec{n}^2}{2mL^2} \equiv \frac{\pi^2 \hbar^2}{2mL^2} n(E)^2$$



$$n(E) = \sqrt{\frac{2mL^2 E}{\pi^2 \hbar^2}}$$

$$n(E)^3 = \left(\frac{2mE}{\pi^2 \hbar^2} \right)^{3/2} V$$

$$\frac{N(E)}{V} = \frac{1}{6\pi^2} \left(\frac{2mE}{\hbar^2} \right)^{3/2} = \frac{k^3}{6\pi^2}$$

where $E \equiv \frac{\hbar^2 k^2}{2m} = \frac{p^2}{2m}$

$$\frac{N(E)}{V} = \frac{P^3}{6\pi^2 \hbar^3} = \frac{4\pi}{3} \frac{P^3}{h^3}$$

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$$N(E) = \int d^3x \frac{4\pi}{3} \frac{P^3}{h^3}$$

$$N(E) = \frac{1}{h^3} \int d^3x \int d^3P$$

$$= \frac{\text{Volume in phase space}}{h^3}$$

⇒ Statistical mechanics