

Electron in a uniform static
magnetic field

$$H = \frac{1}{2m} \left(\vec{p} - \frac{q}{c} \vec{A} \right)^2 + qV$$

For static fields, a convenient gauge is the Coulomb gauge

$$\nabla \cdot \vec{A} = 0$$

Recall $[f(x), p_x] = i\hbar \frac{df}{dx}$

$$\Rightarrow \vec{A} \cdot \vec{p} - \vec{p} \cdot \vec{A} = i\hbar \nabla \cdot \vec{A} = 0$$

The vector potential commutes with momentum in the Coulomb gauge.

Landau levels

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While the quadratic term in \vec{B} can be neglected for bounded systems, such as atoms, in weak fields, this is not true for free particles.

Consider a two-dimensional electron gas in the x - y plane in a field $\vec{B} = B \hat{z}$.

Instead of the symmetric gauge, it is convenient, for this problem to work in the Landau gauge

$$\vec{A} = (-yB, 0, 0)$$

$$H = \frac{1}{2m} \left(p_x - \frac{eBy}{c} \right)^2 + \frac{p_y^2}{2m}$$

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$[p_x, H] = 0$, so seek sol'n of

the form $\Psi(x, y) = e^{ikx} \chi(y)$.

$$H \Psi(x, y) = \frac{1}{2m} \left(\hbar k - \frac{eBy}{c} \right)^2 \Psi + \frac{p_y^2}{2m} \Psi$$

$$E \chi(y) = \left[\frac{p_y^2}{2m} + \frac{1}{2m} \left(\hbar k - \frac{eBy}{c} \right)^2 \right] \chi$$

$$E \chi(y) = \frac{p_y^2}{2m} \chi + \frac{e^2 B^2}{2mc^2} \left(y - \frac{\hbar ck}{eB} \right)^2 \chi$$

$$E \chi = \left[\frac{p_y^2}{2m} + \frac{m \Omega^2}{2} (y - y_0)^2 \right] \chi$$

$$\Omega = \frac{eB}{mc} = \text{cyclotron frequency}$$

$$y_0 = \frac{\hbar ck}{eB}$$

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\Rightarrow shifted harmonic oscillator

$$E_n = \hbar \Omega \left(n + \frac{1}{2} \right),$$

$$n = 0, 1, 2, \dots, \infty$$

Landau levels

Degeneracy

Suppose system has dimensions

L_x, L_y . Impose periodic

BCs in x-direction.

$$e^{ikL_x} = 1 \Rightarrow k = \frac{2\pi n_x}{L_x}$$

Moreover, $0 \leq y_0 \leq L_y$

$$0 \leq \frac{\hbar c}{eB} \frac{2\pi n_x}{L_x} \leq L_y$$

$$0 \leq n_x \leq \frac{eB}{\hbar c} L_x L_y$$

Total number of states in a given Landau level of energy $E_n = \hbar\Omega(n + 1/2)$ is

thus
$$N_n = \frac{eB}{\hbar c} A$$

$$N_n = \frac{BA}{\phi_0}$$

$$\phi_0 = \frac{hc}{e}$$

(10)

(normal flux quantum)

Filling factor ν

For a total # N of electrons,
the total # ν of Landau
levels populated is

$$\nu = \frac{N}{N_n} = \frac{N \phi_0}{BA}$$

$$= \frac{\# \text{ electrons}}{\text{flux quantum}}$$

Quantum Hall Effect

$$\frac{h}{e^2} = 25.8 \text{ k}\Omega$$

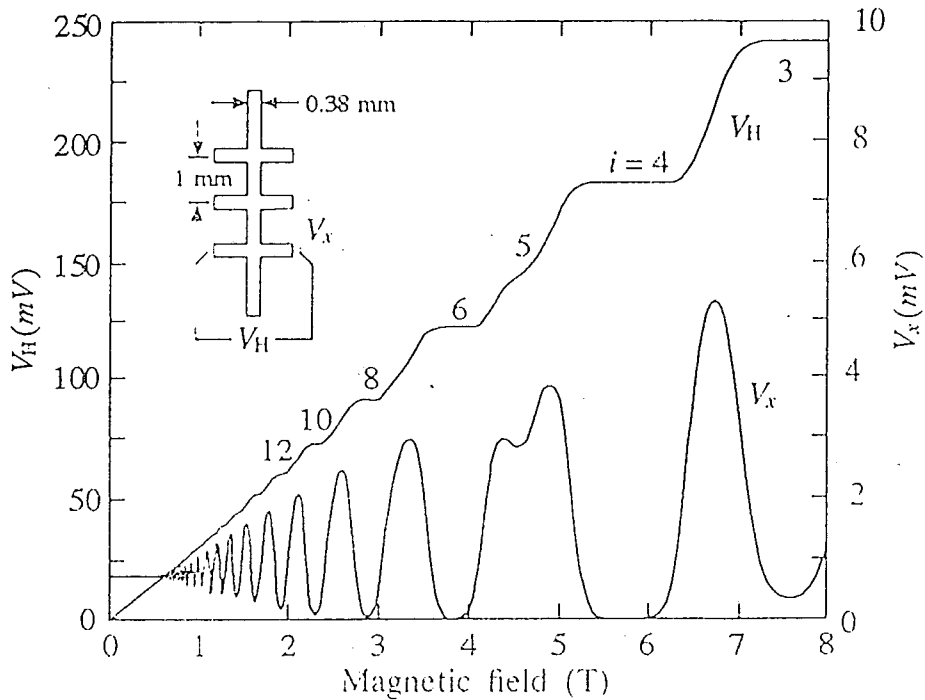


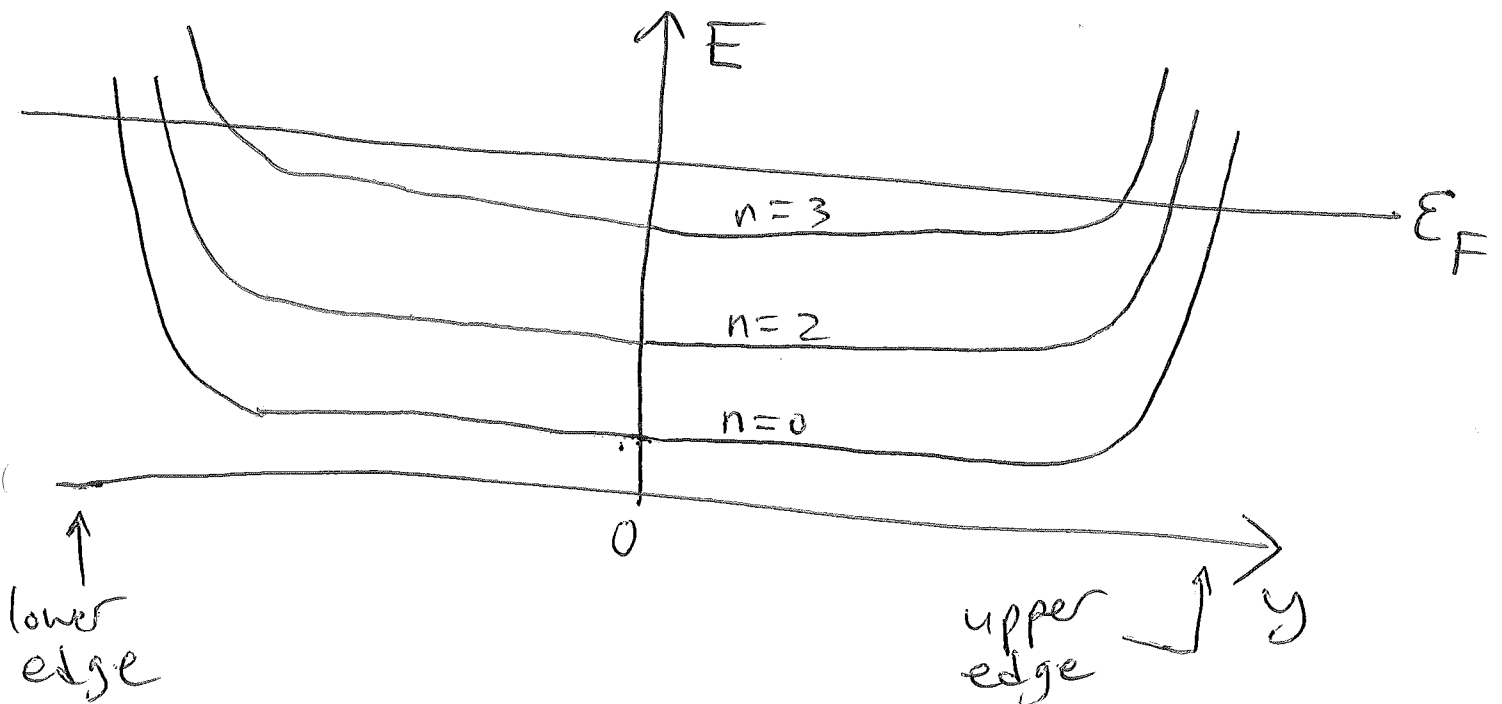
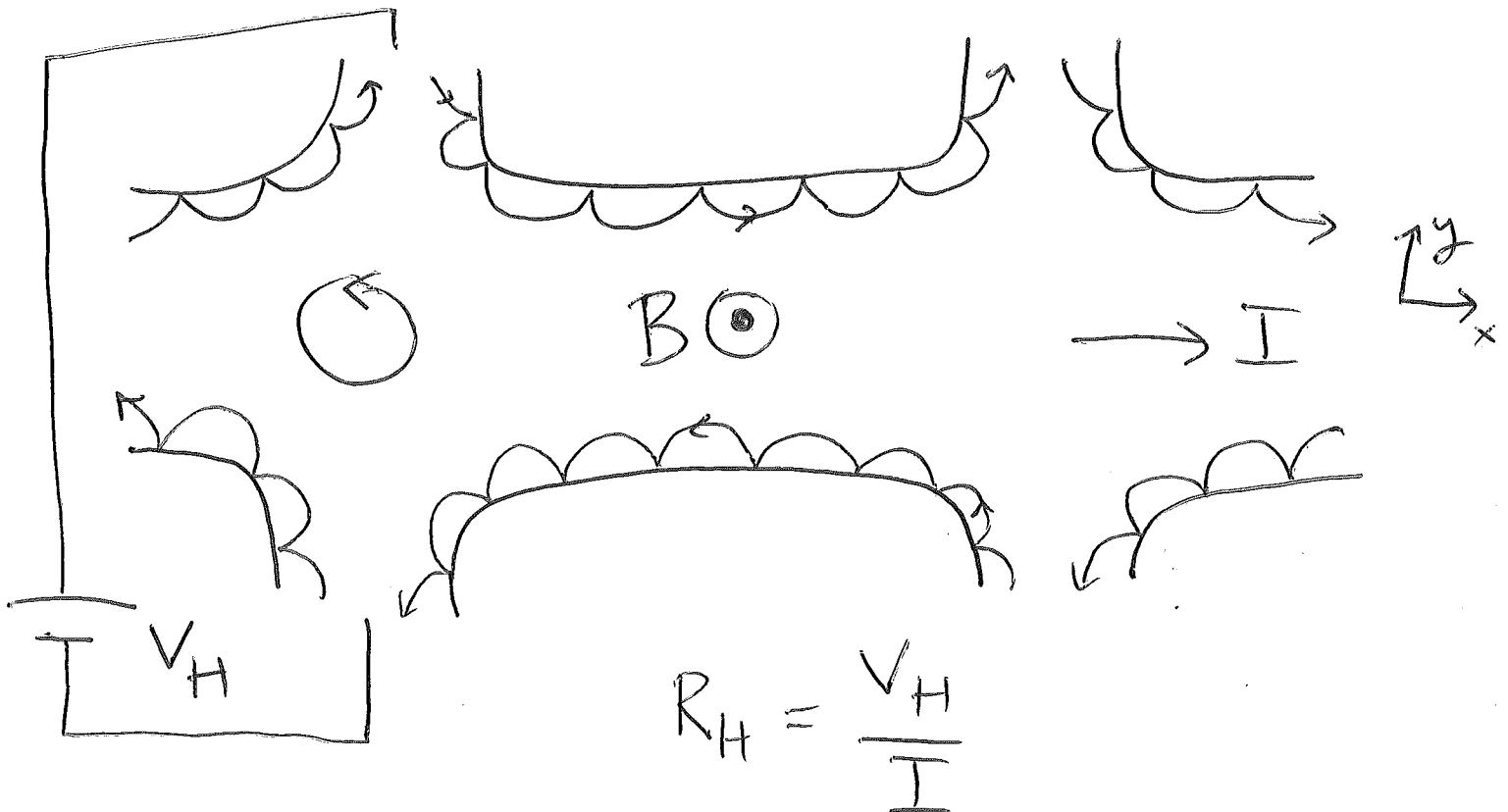
Fig. 1.4.2. Measured longitudinal and transverse voltages for a modulation-doped GaAs film at $T = 1.2 \text{ K}$ ($I = 25.5 \mu\text{A}$). Reproduced with permission from Fig. 1 of M. E. Cage, R. F. Dziuba and B. F. Field (1985), *IEEE Trans. Instrum. Meas.* IM-34, 301. © 1985 IEEE

$$R_H = \frac{h}{4e^2}$$

classically: $R_H = \frac{B}{ne}$

Quantum Hall effect

Current carried by quantized "edge states":



Group velocity

$$V_x^{(n)} = \frac{1}{\hbar} \frac{\partial E_n}{\partial k_x}$$

Recall $y_0 = \frac{\hbar c k_x}{eB}$

$$V_x^{(n)} = \frac{1}{\hbar} \frac{\hbar c}{eB} \frac{\partial E_n}{\partial y_0}$$

$$E_n(y_0) \approx \hbar \Omega \left(n + \frac{1}{2} \right) + U(y_0)$$

$U(y) =$ confining potential

$$V_x^{(n)} \approx \frac{c}{eB} \frac{\partial U}{\partial y_0}$$

approx.
indep.
of quantum #
 \hbar

States at upper edge travel right, lower go left.

These edge states form ideal one-dimensional channels. The # of channels is equal to the # of Landau levels lying below $\epsilon_F \Rightarrow \nu$ channels.

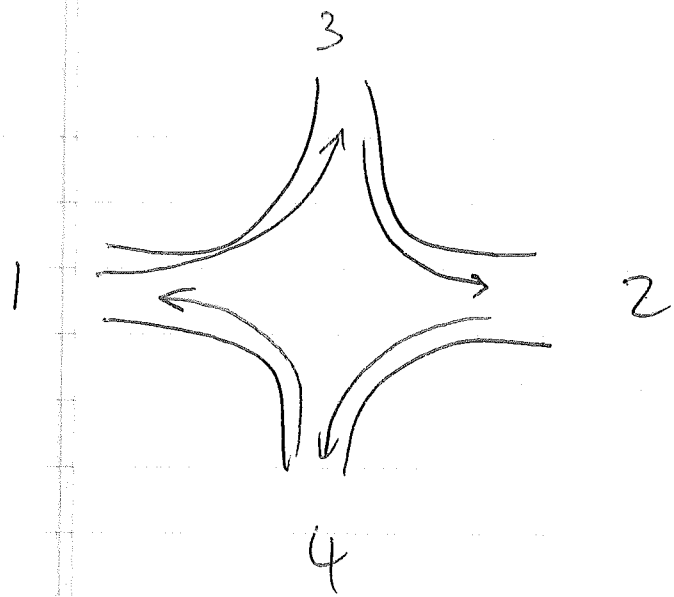
$$R_H = \frac{V_H}{I} = \frac{h}{\nu e^2}$$

Recall degeneracy $N_n = \frac{eBA}{hc}$

$$\nu = \text{Int} \left\{ \frac{N}{N_n} \right\} = \text{Int} \left\{ \int_{2D} \frac{hc}{eB} \right\}$$

Classical limit $\nu \approx \int_{2D} \frac{hc}{eB}$

$$R_H \approx \frac{h}{\nu e^2} \approx \frac{B}{\int_{2D} ec} \quad \checkmark$$



$$I_4 = \frac{ve}{h} (\mu_2 - \mu_4) = 0 \quad (\text{voltage probe})$$

$$\Rightarrow \mu_4 = \mu_2$$

$$I_3 = \frac{ve}{h} (\mu_1 - \mu_3) = 0 \quad (\text{voltage probe})$$

$$\Rightarrow \mu_3 = \mu_1$$

$$I_2 = \frac{ve}{h} (\mu_3 - \mu_2) = \frac{ve}{h} (\mu_1 - \mu_2) = eV_H$$