

## Sommerfeld Expansion

For fermionic systems, we often need to compute averages of the form:

$$\langle A \rangle = \int_{-\infty}^{\infty} d\varepsilon A(\varepsilon) f(\varepsilon),$$

$$\text{where } f(\varepsilon) = \frac{1}{e^{\beta(\varepsilon - \mu)} + 1}.$$

Assuming  $A(-\infty) = 0$ , one can integrate by parts:

$$\langle A \rangle = \int_{-\infty}^{\infty} d\varepsilon \left[ \int_{-\infty}^{\varepsilon} d\varepsilon' A(\varepsilon') \right] \left( -\frac{\partial f}{\partial \varepsilon} \right)$$

$-\frac{\partial F}{\partial \epsilon}$  is sharply peaked about  $\mu$  (2)

$\epsilon = \mu$  and has unit integral.

Thus we can expand

$$\int_{-\infty}^{\epsilon} d\epsilon' A(\epsilon') = \int_{-\infty}^{\mu} d\epsilon' A(\epsilon') + \int_{\mu}^{\epsilon} d\epsilon' A(\epsilon')$$

Using  $A(\epsilon') = A(\mu) + A'(\mu)(\epsilon' - \mu) + \dots$

$$= \sum_{n=0}^{\infty} \left. \frac{d^n A}{d\epsilon'^n} \right|_{\mu} \frac{(\epsilon' - \mu)^n}{n!}$$

$$\int_{\mu}^{\epsilon} d\epsilon' A(\epsilon') = A(\mu)(\epsilon - \mu) + A'(\mu) \frac{(\epsilon - \mu)^2}{2} + \dots$$

$$= \sum_{n=0}^{\infty} \frac{d^n A(\mu)}{d\mu^n} \frac{(\epsilon - \mu)^{n+1}}{(n+1)!}$$

$-\frac{\partial f}{\partial \epsilon}$  is an even function of  $\epsilon$  3

$\epsilon - \mu$ , so only the even powers in the series contribute to  $\langle A \rangle$ .

$$\langle A \rangle = \int_{-\infty}^{\mu} d\epsilon A(\epsilon) + \sum_{n=1}^{\infty} \frac{d^{2n-1} A(\mu)}{d\mu^{2n-1}} \frac{1}{(2n)!}$$

$$\times \int_{-\infty}^{\infty} d\epsilon (\epsilon - \mu)^{2n} \left( -\frac{\partial f}{\partial \epsilon} \right)$$

let  $x = \beta(\epsilon - \mu)$

$$\langle A \rangle = \int_{-\infty}^{\mu} d\epsilon A(\epsilon) + \sum_{n=1}^{\infty} a_n (k_B T)^{2n} \frac{d^{2n-1} A(\mu)}{d\mu^{2n-1}}$$

where

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$$a_n = \frac{1}{(2n)!} \int_{-\infty}^{\infty} dx \frac{x^{2n}}{\left(e^{\frac{x}{2}} + e^{-\frac{x}{2}}\right)^2}$$

$$a_1 = \frac{\pi^2}{6}, \quad a_2 = \frac{7\pi^4}{360}, \quad \text{etc.}$$

So

$$\langle A \rangle = \int_{-\infty}^{\mu} d\varepsilon A(\varepsilon) + \frac{\pi^2}{6} (k_B T)^2 A'(\mu) + \frac{7\pi^4}{360} (k_B T)^4 A'''(\mu) + \dots$$

Example: Specific heat of a Fermi gas (revisited).

$$E = \int_{-\infty}^{\infty} d\varepsilon \varepsilon D(\varepsilon) f(\varepsilon)$$

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$$A(\varepsilon) = \varepsilon D(\varepsilon)$$

$$E = \int_{-\infty}^{\mu} d\varepsilon \varepsilon D(\varepsilon) + \frac{\pi^2}{6} (k_B T)^2 \frac{d(\mu D(\mu))}{d\mu} + \mathcal{O}(k_B T)^4$$

How does  $\mu$  depend on temperature?

$$\left. \frac{\partial \mu}{\partial T} \right|_{N, V} = - \frac{\left. \frac{\partial N}{\partial T} \right|_{\mu, V}}{\left. \frac{\partial N}{\partial \mu} \right|_{T, V}}$$

$$N = \int_{-\infty}^{\infty} d\varepsilon f(\varepsilon) D(\varepsilon)$$

$$= \int_0^{\mu} d\varepsilon D(\varepsilon) + \frac{\pi^2}{6} (k_B T)^2 D'(\mu)$$

$$\Rightarrow \left. \frac{\partial \mu}{\partial T} \right|_{N, V} = -\frac{\pi^2}{3} k_B^2 T \frac{D'(\mu)}{D(\mu)} + \mathcal{O}(k_B T)^3$$

$$\mu = \varepsilon_F - \frac{\pi^2}{6} (k_B T)^2 \frac{D'(\varepsilon_F)}{D(\varepsilon_F)}$$

$$D(\varepsilon) = \frac{V}{2\pi^2} \left( \frac{2m}{\hbar^2} \right)^{3/2} \varepsilon^{1/2} \quad \text{in 3D}$$

$$\frac{D'(\varepsilon_F)}{D(\varepsilon_F)} = \frac{1}{2\varepsilon_F}$$

$$\mu \simeq \varepsilon_F - \frac{\pi^2}{12} \frac{(k_B T)^2}{\varepsilon_F}$$

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$$E = \int_0^{\varepsilon_F} d\varepsilon \varepsilon D(\varepsilon) + \frac{\pi^2}{6} (k_B T)^2 D(\varepsilon_F)$$

$$+ \varepsilon_F \left\{ (\mu - \varepsilon_F) D(\varepsilon_F) + \frac{\pi^2}{6} (k_B T)^2 D(\varepsilon_F) \right\}$$

→ 0

$$\Rightarrow C_V = \frac{\pi^2}{3} k_B^2 T D(\varepsilon_F) \quad \checkmark$$