

$$1) \quad \langle E \rangle = \frac{1}{2m} \langle p_x^2 \rangle + \frac{m\omega^2}{2} \langle x^2 \rangle$$

$$(\Delta x)^2 = \langle (x - \langle x \rangle)^2 \rangle = \langle x^2 \rangle - \langle x \rangle^2 \leq \langle x^2 \rangle$$

$$(\Delta p_x)^2 \leq \langle p_x^2 \rangle$$

$$\Rightarrow \langle E \rangle \geq \frac{1}{2m} (\Delta p_x)^2 + \frac{m\omega^2}{2} (\Delta x)^2$$

$$\text{Also } (\Delta x)^2 (\Delta p_x)^2 \geq \left(\frac{\hbar}{2}\right)^2$$

$$\langle E \rangle \geq \frac{1}{2m} \frac{\hbar^2}{4(\Delta x)^2} + \frac{m\omega^2}{2} (\Delta x)^2$$

minimize over Δx :

$$0 = \frac{\partial \langle E \rangle}{\partial \Delta x} = -\frac{1}{2m} \frac{\hbar^2}{2(\Delta x)^3} + m\omega^2 \Delta x$$

$$(\Delta x)^4 = \frac{\hbar^2}{4m^2\omega^2} \quad (\Delta x)^2 = \frac{\hbar}{2m\omega}$$

$$\langle E \rangle \geq \frac{\hbar^2}{8m} \frac{2m\omega}{\hbar} + \frac{m\omega^2}{2} \frac{\hbar}{2m\omega} = \frac{\hbar\omega}{2}$$

This is the ground-state energy.

$$2) \text{ Let } |\psi\rangle = \alpha|1\rangle + \beta|2\rangle$$

$$\hat{H}|\psi\rangle = E|\psi\rangle$$

$$\langle 1|\hat{H}|\psi\rangle = E\langle 1|\psi\rangle$$

$$\langle 2|\hat{H}|\psi\rangle = E\langle 2|\psi\rangle$$

$$\begin{pmatrix} H_{11} & H_{12} \\ H_{12} & H_{22} \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = E \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

$$0 = \det \begin{pmatrix} H_{11} - E & H_{12} \\ H_{12} & H_{22} - E \end{pmatrix} = (E - H_{11})(E - H_{22}) - H_{12}^2$$

$$0 = E^2 - (H_{11} + H_{22})E - H_{12}^2 + H_{11}H_{22}$$

$$E_{\pm} = \frac{H_{11} + H_{22}}{2} \pm \sqrt{\left(\frac{H_{11} + H_{22}}{2}\right)^2 + H_{12}^2 - H_{11}H_{22}}$$

$$E_{\pm} = \frac{H_{11} + H_{22}}{2} \pm \sqrt{\left(\frac{H_{11} - H_{22}}{2}\right)^2 + H_{12}^2}$$

Let $\bar{E} = \frac{H_{11} + H_{22}}{2}$, $\frac{H_{22} - H_{11}}{2} = \Delta$,

$$\frac{\Delta}{H_{12}} = x, \quad H_{12} = V.$$

$$E_{\pm} = \bar{E} \pm |V| \sqrt{1 + x^2}$$

$$(H_{11} - E_{\pm}) \alpha_{\pm} + H_{12} \beta_{\pm} = 0$$

$$\beta_{\pm} = \frac{E_{\pm} - H_{11}}{H_{12}} \alpha_{\pm}$$

$$\beta_{\pm} = \frac{\Delta \pm |V| \sqrt{1 + x^2}}{V} \alpha_{\pm} = (x \pm \sqrt{1 + x^2}) \alpha_{\pm}$$

Eigenfunctions:

$$|\psi_{\pm}\rangle = \alpha_{\pm} (|1\rangle + (x \pm \sqrt{1 + x^2}) |2\rangle)$$

Normalization:

$$1 = \langle \psi_{\pm} | \psi_{\pm} \rangle = |\alpha_{\pm}|^2 (1 + (x \pm \sqrt{1+x^2})^2)$$

$$\alpha_{\pm} = \frac{1}{\sqrt{1 + (x \pm \sqrt{1+x^2})^2}}$$

3) Since $\{|a', b'\rangle\}$ forms a complete basis in Hilbert space, a general state can be written

$$|\psi\rangle = \sum_{a', b'} C_{a'b'} |a'b'\rangle$$

$$[\hat{A}, \hat{B}] |\psi\rangle = \sum_{a', b'} C_{a'b'} [\hat{A}, \hat{B}] |a'b'\rangle$$

$$\text{But } [\hat{A}, \hat{B}] |a'b'\rangle = (a'b' - b'a') |a'b'\rangle = 0$$

$$\text{So } [\hat{A}, \hat{B}] |\psi\rangle = 0 \quad \forall |\psi\rangle \in \mathcal{H}$$

$$\Rightarrow [\hat{A}, \hat{B}] = 0$$

4) Suppose $A|\psi\rangle = a|\psi\rangle$ and $B|\psi\rangle = b|\psi\rangle$

$$\{A, B\} |\psi\rangle = (AB + BA) |\psi\rangle = 2ab |\psi\rangle = 0$$

This is possible if $a = 0$ or $b = 0$.

$$5) \quad [A_1, A_2] \neq 0, \quad [A_1, H] = 0, \quad [A_2, H] = 0.$$

Suppose $H|v\rangle = E_v|v\rangle$ with E_v distinct eigenvalues (no degeneracy). Then, from the theorem proved at the beginning of lecture 12, $|v\rangle$ must also be an eigenstate of A_1 and A_2 .

$$A_1|v\rangle = a_v|v\rangle, \quad A_2|v\rangle = b_v|v\rangle.$$

$$[A_1, A_2]|v\rangle = (a_v b_v - b_v a_v)|v\rangle = 0. \quad \text{But}$$

a general state can be written

$$|\psi\rangle = \sum_v c_v |v\rangle.$$

$$[A_1, A_2]|\psi\rangle = \sum_v c_v [A_1, A_2]|v\rangle = 0.$$

This is a contradiction, since $[A_1, A_2] \neq 0$.
 \therefore The energy eigenstates must in general be degenerate.

An exception can occur if $a_v = 0$ and/or $b_v = 0$. For instance, the $l=0$ states in a central potential have $L_x = L_y = L_z = 0$ even though $[L_x, L_y] \neq 0$, etc.
