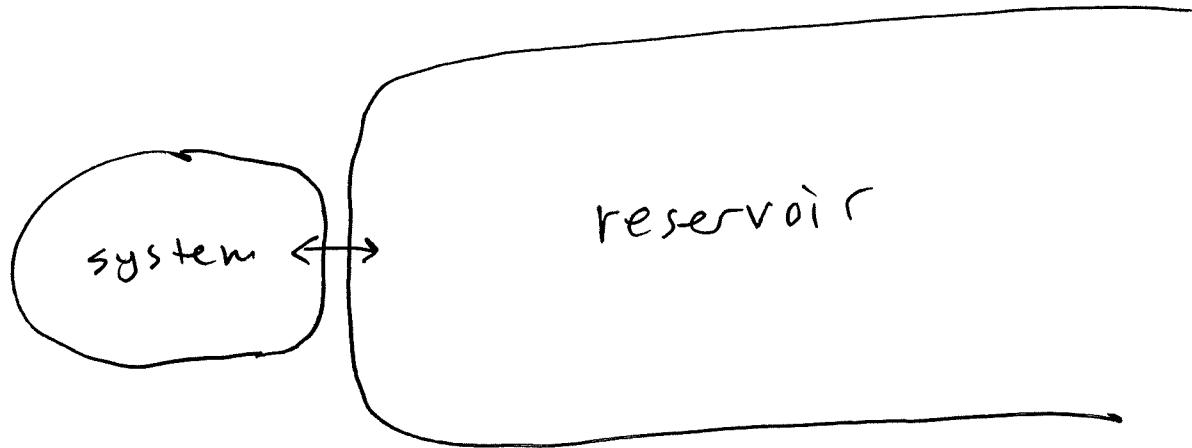


Lec. 12

Energy - time uncertainty

Open quantum system



$$\hat{H} = \hat{H}_{\text{sys}} + \hat{H}_{\text{res}} + \hat{H}_{\text{int}}$$

$$\hat{H}_{\text{sys}} = \sum_{n,m} H_{nm} |n\rangle\langle m|$$

$$\hat{H}_{\text{res}} = \sum_k \varepsilon_k |k\rangle\langle k|$$

$$\hat{H}_{\text{int}} = \sum_{n,k} (V_{nk} |n\rangle\langle k| + \text{H.c.})$$

Retarded Green's function

Let $G^r(t) = -i \frac{\theta(t)}{\hbar} \hat{U}(t)$

$$\hat{U}(t) = \exp\left(-i\frac{\hat{H}t}{\hbar}\right)$$

$$i\hbar \frac{d}{dt} G^r(t) = S(t) + \hat{H} G^r(t)$$

$$\left(i\hbar \frac{d}{dt} - \hat{H}\right) G^r(t) = S(t)$$

Suppose we are only interested in/ able to measure the observables in the small quantum system, not the large reservoir weakly coupled to it.

These can be obtained from the submatrix of $G^r(t)$:

$$G_{nm}^r(t) = -i\frac{\theta(t)}{\hbar} \langle n | \hat{U}(t) | m \rangle,$$

where $|n\rangle, |m\rangle \in \mathcal{H}_{sys} = \text{Hilbert space}^{\text{system}}$

$$i\hbar \frac{d}{dt} G_{nm}^r(t) = S(t) \langle n | m \rangle - i\frac{\theta(t)}{\hbar} \langle n | \hat{H} \hat{U}(t) | m \rangle$$

Fourier transform of $G^r(+)$

$$G^r(\omega) = \int_{-\infty}^{\infty} dt e^{i\omega t} G^r(+)$$

$$(\hbar\omega - \hat{H}) G^r(\omega) = \underline{1}$$

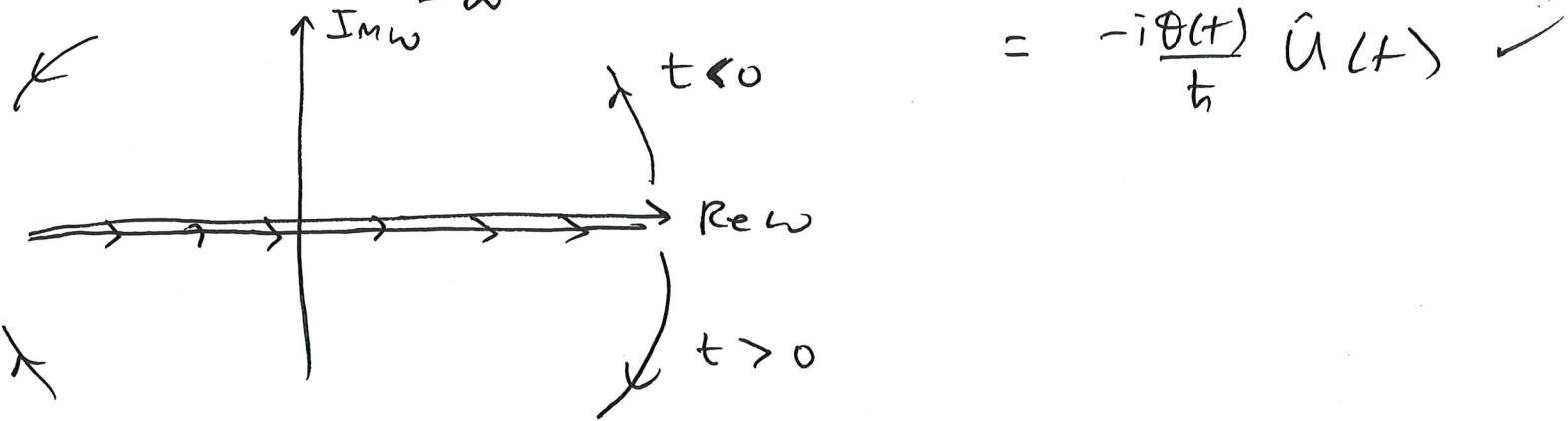
$$G^r(\omega) = (\hbar\omega - \hat{H} + i\sigma^+)^{-1}$$

\underline{C} needed for causality

$$\text{Let } \hat{H}|v\rangle = E_v |v\rangle$$

$$G^r(\omega) = \sum_v \frac{|v\rangle\langle v|}{\hbar\omega - E_v + i\sigma^+}$$

$$G^r(t) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} G^r(\omega) = -i \frac{\theta(t)}{\hbar} \sum_v e^{-i\frac{E_v}{\hbar} t} |v\rangle\langle v|$$



Spectral function

$$A(\omega) = \frac{i}{2\pi} [G^r(\omega) - G^a(\omega)] \\ = \sum_v |v\rangle\langle v| \delta(\hbar\omega - E_v)$$

Density of states

$$g(\omega) = \text{Tr} \{ A(\omega) \} = \sum_v \delta(\hbar\omega - E_v)$$

→ spectrum of system

Here we have used the identity

$$\lim_{\eta \rightarrow 0^+} \frac{1}{x + i\eta} = P \frac{1}{x} - i\pi \delta(x)$$

(see books on complex variable theory.)

$$G^a(\omega) = [G^r(\omega)]^+$$

$$i\hbar \frac{d}{dt} G_{nm}^r(t) = \delta(t) \delta_{nm}$$

$$-i\frac{\theta(t)}{\hbar} \left(\sum_{n'} \langle n | \hat{H} | n' \rangle \langle n' | \hat{U}(t) | m \rangle + \sum_k \langle n | \hat{H} | k \rangle \langle k | \hat{U}(t) | m \rangle \right)$$

$$i\hbar \frac{d}{dt} G_{nm}^r(t) = \delta(t) \delta_{nm} + \sum_{n'} H_{nn'} G_{n'm}^r(t)$$

$$+ \sum_k V_{nk} G_{km}^r(t)$$

$$i\hbar \frac{d}{dt} G_{km}^r(t) = -i\frac{\theta(t)}{\hbar} \langle k | \hat{H} \hat{U}(t) | m \rangle$$

$$= -i\frac{\theta(t)}{\hbar} \left(\sum_{k'} \langle k | \hat{H} | k' \rangle \langle k' | \hat{U}(t) | m \rangle + \sum_{n'} \langle k | \hat{H} | n' \rangle \langle n' | \hat{U}(t) | m \rangle \right)$$

$$i\hbar \frac{d}{dt} G_{km}^r(t) = \epsilon_k G_{km}^r(t) + \sum_{n'} V_{kn} G_{n'm}^r(t)$$

Fourier transform

$$(h\omega - \varepsilon_k + i0^+) G_{km}^r(\omega) = \sum_{n'} V_{kn'} G_{n'm}^r(\omega)$$

↑
causality

$$G_{km}^r(\omega) = \frac{1}{h\omega - \varepsilon_k + i0^+} \sum_{n'} V_{kn'} G_{n'm}^r(\omega)$$

$$\sum_{n'} (h\omega \delta_{nn'} - H_{nn'} + i0^+) G_{n'm}^r(\omega) = S_{nm}$$

$$+ \sum_k V_{nk} G_{km}^r(\omega)$$

$$= S_{nm} + \sum_{k, n'} \frac{V_{nk} V_{kn'}}{h\omega - \varepsilon_k + i0^+} G_{n'm}^r(\omega)$$

$$\text{Let } \sum_{nn'}^r(\omega) \equiv \sum_k \frac{V_{nk} V_{kn'}}{h\omega - \varepsilon_k + i0^+}$$

$$(1 h\omega - \hat{H}_{\text{sys}} - \sum^r(\omega)) G_{\text{sys}}^r(\omega) = 1$$

$$G_{\text{sys}}^r(\omega) = (1 h\omega - \hat{H}_{\text{sys}} - \sum^r(\omega))^{-1}$$

$$\Sigma(\omega) = \sigma(\omega) - \frac{i}{2} \Gamma(\omega),$$

where

$$\sigma_{nm}(\omega) = \langle \Phi | \sum_k \frac{V_{nk} V_{km}}{\hbar\omega - \epsilon_k}$$

$$\Gamma_{nm}(\omega) = 2\pi \sum_k V_{nk} V_{km} \delta(\hbar\omega - \epsilon_k)$$

$$= 2\pi \overline{V_{nk} V_{km}} g_{\text{res}}(\hbar\omega)$$

(density of states
of reservoir)

cf. Fermi's golden rule

$\sigma(\omega)$ simply shifts the system energy levels, but $\Gamma(\omega)$ moves the energy values off the real axis, leading to broadening, and is qualitatively more significant.

Simple case $\sigma(\omega) = 0$

$$\Gamma(\omega) = \Gamma \perp$$

$$G_{sys}^r(\omega) = \left(\hat{H}_{sys} - i\Gamma/2 \right)^{-1}$$

Let $\hat{H}_{sys} |\alpha\rangle = E_\alpha |\alpha\rangle$.

$$G_{sys}^r(\omega) = \sum_{\alpha} \frac{|\alpha\rangle\langle\alpha|}{\hbar\omega - E_{\alpha} + i\Gamma/2}$$

$$\begin{aligned} A_{sys}(\omega) &= \frac{i}{2\pi} \left[G_{sys}^r(\omega) - G_{sys}^s(\omega) \right] \\ &= \frac{i}{2\pi} \sum_{\alpha} |\alpha\rangle\langle\alpha| \left(\frac{1}{\hbar\omega - E_{\alpha} + i\Gamma/2} - \frac{1}{\hbar\omega - E_{\alpha} - i\Gamma/2} \right) \\ &= \sum_{\alpha} \frac{|\alpha\rangle\langle\alpha| \Gamma/2\pi}{(\hbar\omega - E_{\alpha})^2 + (\Gamma/2)^2} \end{aligned}$$

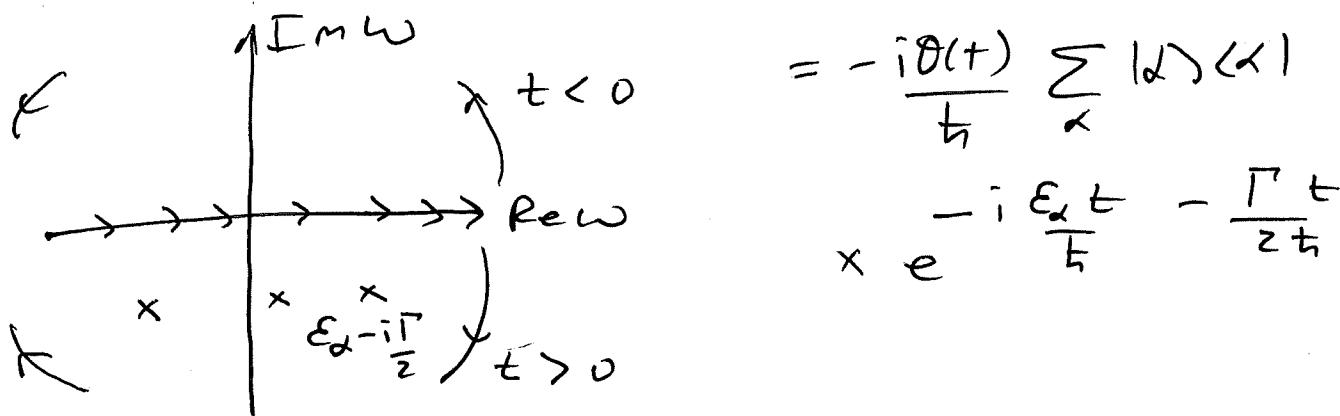
$$\begin{aligned} g_{sys}(\omega) &= \text{Tr} \{ A_{sys}(\omega) \} = \text{Tr}_{sys} \{ A(\omega) \} \\ &= \sum_{\alpha} \frac{\Gamma/2\pi}{(\hbar\omega - E_{\alpha})^2 + (\Gamma/2)^2} \end{aligned}$$

Energy levels broadened into Lorentzians,
 $\text{FWHM} = \Gamma$.

Let's Fourier transform back
to the time domain:

$$G^r(t) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} G^r(\omega) e^{-i\omega t}$$

$$= \sum_{\alpha} |\alpha\rangle \langle \alpha| \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{e^{-i\omega t}}{\hbar\omega - \epsilon_{\alpha} + i\Gamma/2}$$



$$|G_{\alpha\alpha}^r(t)\rangle^2 = \frac{1}{\hbar^2} \theta(t) e^{-\frac{\Gamma}{\hbar} t}$$

$$\equiv \frac{\theta(t)}{\hbar^2} e^{-t/\tau}$$

lifetime $\tau = \frac{\hbar}{\Gamma} \equiv \Delta t$

$$\Delta E = \frac{\Gamma}{2} = \text{HWHM}$$

$$\Delta E \Delta t = \frac{\Gamma}{2} \frac{\hbar}{\Gamma} = \frac{\hbar}{2}$$

Energy-time
uncertainty
relation