

# Charge Symmetry Breaking in $dd \rightarrow {}^4\text{He} \pi^0$

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- Introduction
- The Indiana  $dd \rightarrow {}^4\text{He} \pi^0$  experiment.
- Plane wave calculation including heavy meson exchange and  $\pi$ - $\eta$  meson mixing.
- TRIUMF CSB experiment  $np \rightarrow d \pi^0$ 
  - Anders Gardestig and C. J. Horowitz

# Isospin Violation Probes Strongly Interacting Systems

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- Determine  $m_u, m_d$ : fundamental and poorly known parameters of Standard Model.
- Determine strong interaction phases and amplitudes for CP violation studies.
  - Need two interfering amplitudes for ~~CP~~.  
 $B \rightarrow \rho\rho$  can interfere with  $B \rightarrow \rho\omega \rightarrow \rho\rho$  because of  $\rho\omega$  mixing.
  - Extracted ~~CP~~ parameters such as  $\sin 2\alpha$  from other decays can be impacted by isospin violation [S. Gardner].

# Determine Weak Interaction Matrix Elements for Standard Model Tests

- Super-allowed  $\beta$  decays give quark mixing  $V_{ud}$  after correcting for isospin violation.

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1 ?$$

Brookhaven E865 result for  $K^+ \rightarrow \pi^0 e^+ \nu$  branching ratio changes  $V_{us}$  and (if true) makes matrix unitary.

- NUTEV Anomaly Sensitive to CSB

$$(\langle \sigma_{NC}^{\nu} \rangle - \langle \sigma_{NC}^{\bar{\nu}} \rangle) / (\langle \sigma_{CC}^{\nu} \rangle - \langle \sigma_{CC}^{\bar{\nu}} \rangle) = 1/2 - 2\sin^2\theta_W.$$

- Assumes charge symmetry:  $p \rightarrow n$  and  $u \rightarrow d$
- Londergan and Thomas find that CSB in  $u_{\nu}(x)$  in  $p$  compared to  $d_{\nu}(x)$  in  $n$  related to  $np$  and  $ud$  mass differences can explain  $\approx 1/3$  of anomaly.

# Separate E+M and Quark Mass Contributions to Isospin Violation

- Two apparently independent sources:
  - E+M makes proton heavier than neutron.
  - $\Delta m \equiv m_d - m_u$  makes n heavier than p.
- $\Delta M \equiv M_n - M_p = 1.29 \text{ MeV}$  one of most basic and important parameters in nuclear astrophysics.
  - If  $M_p \approx M_n$  Big Bang makes mostly  ${}^4\text{He}$ .
  - If  $M_p > M_n$  Big Bang makes  ${}^4\text{He}$  and n (stable). H unstable.
  - In either case Universe drastically different.
- How to separate E+M and  $\Delta m$  effects?
  - Use different density dependence.
  - Study isospin violation in  $\pi^0$ -N scattering.

# In a high density system

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- Kinematics suppresses quark mass effects by large Fermi momentum.  $(k_F^2+m^2)^{1/2} \sim k_F + m^2/2k_F$
- E+M effects grow as quarks move closer together. For example Coulomb exchange energy  $\sim \alpha k_F$ .
- In high density limit, isospin violation is dominated by E+M with quark mass terms unimportant. In this limit the proton is heavier than neutron.

[CJH+J. Piekarewicz, PRC**63**, (2001) 011303]

# Isospin Violating $\pi^0$ -N Scattering

## (U. van Kolck)

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- Chiral sym relates nucleon mass term to  $\pi$ -N scattering.
- Isospin breaking in nucleon mass term (n-p mass difference) related to isospin violating  $\pi$ -N scattering.
- Isospin viol. with symmetries of quark masses:  
$$L_{qm} = \delta m / 2 [N^\dagger \tau_3 N - 2 N^\dagger \pi_3 \tau \cdot \pi N / F_\pi^2]$$
- With symmetries of hard photon exchange:  
$$L_{em} = \bar{\delta} m / 2 [N^\dagger \tau_3 N + 2 N^\dagger (\pi_3 \tau \cdot \pi - \pi^2 \tau_3) / F_\pi^2]$$
- $\Delta M = \delta m + \bar{\delta} m$  and measuring  $\pi^0$ -N scattering can decompose  $\Delta M$  into quark mass and E+M parts.
- $\pi^0$  decay makes  $\pi^0$  scattering hard. Look at isospin violation in  $\pi^0$  production instead.

# Cooler

# CSB



the search

for  $d+d \rightarrow \alpha\pi^0$

# *The Cooler-CSB Experiment at the Indiana University Cooler Synchrotron: A Test of Charge Symmetry Breaking and Isospin Conservation via the Reaction $d+d \rightarrow \alpha+\pi^0$*

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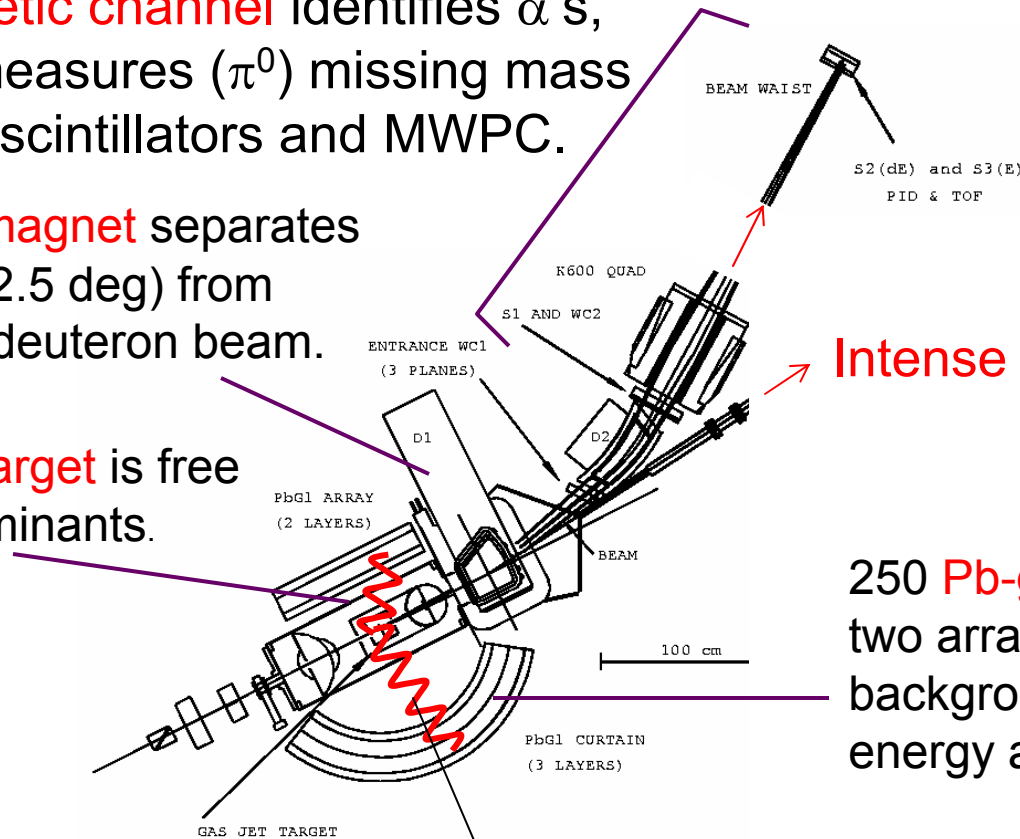


# $d+d \rightarrow {}^4\text{He} + \pi^0$ at the IUCF Cooler

**Magnetic channel** identifies  $\alpha$ 's, and measures ( $\pi^0$ ) missing mass using scintillators and MWPC.

**6 degree magnet** separates  $\alpha$ 's (bent 12.5 deg) from circulating deuteron beam.

**Gas jet target** is free of contaminants.



**Intense deuteron beam**

**250 Pb-glass detectors** in two arrays remove trigger background and check  $\pi^0$  energy and  $\gamma$  coplanarity.

**See both  $\gamma$ s** from  $\pi^0$  decay

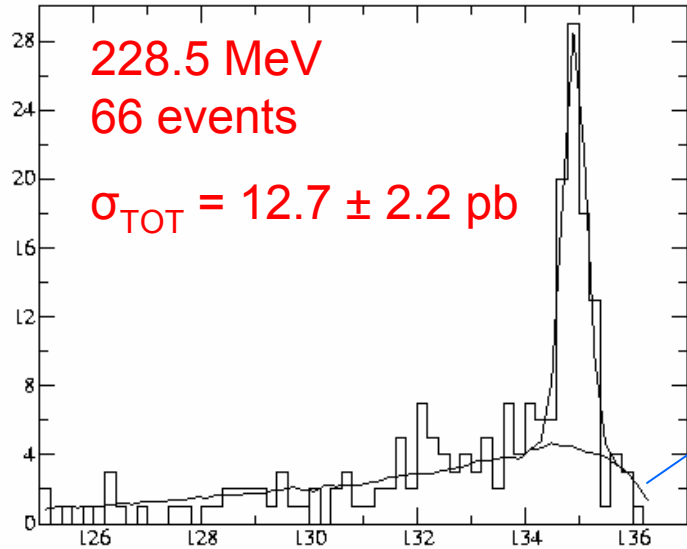
# Experimental Results

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- Data was taken at two near threshold energies: 228.5 MeV ( $\eta_\pi=0.1434$ ) and 231.8 MeV ( $\eta_\pi=0.209$ ). Also some d+d elastic scattering data taken to help check initial state distortions.
- See a clear pion signal reconstructing to pion mass at both energies while background moves with phase space. Low background for two gamma plus He events.
- Luminosity calibration from dd and pd elastic scattering.

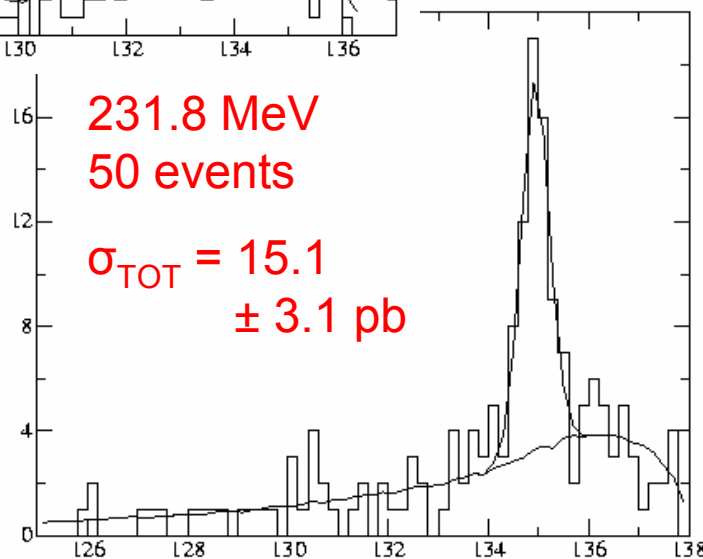
# PRELIMINARY RESULTS

Events in these spectra must satisfy:  
correct pulse height in channel scintillators  
usable wire chamber signals  
good Pb-glass pulse height and timing



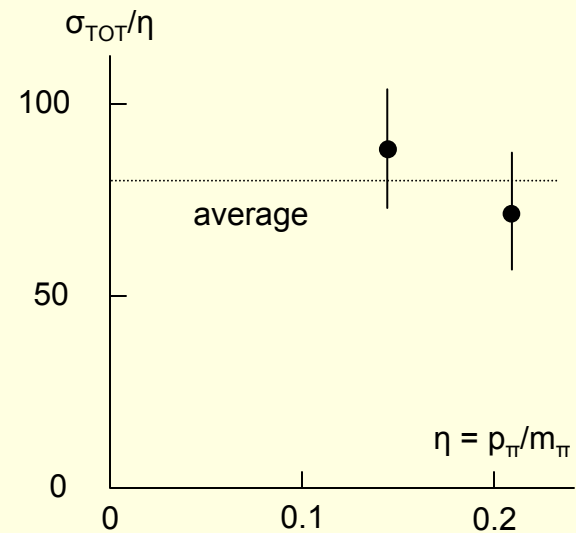
Background shape based on calculated double radiative capture, corrected by empirical channel acceptance using  $^4\text{He}$ .

Spectra are essentially free of random background.



missing mass (MeV)

Cross sections are consistent with S-wave pion production.



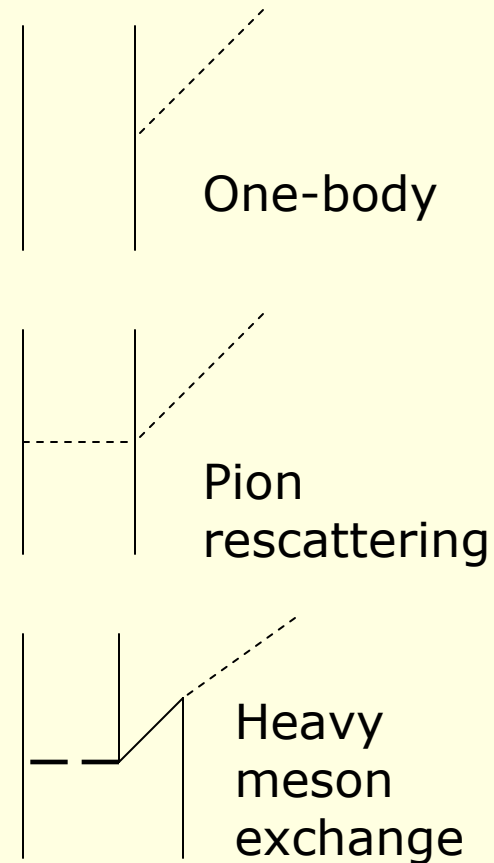
# CSB Theory Collaboration

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- **Antonio Fonseca** (dd scattering wave function), **Anders Gardestig** (plane wave calc.), Chris Hanhart (chiral counting and loops), Chuck Horowitz (CSB operators), Gerry Miller, **Andreas Nogga** ( $^4\text{He}$  wave function), Bira van Kolck, ... have formed collaboration. Four workshops at Seattle, Albuquerque, Bloomington, and Philly.
- There will be a CSB Workshop in Seattle ~ Oct 20, 2003 and all are welcome.
- We have preliminary plane wave results.
- Full distorted wave calculations with realistic  $^4\text{He}$  wave function underway.

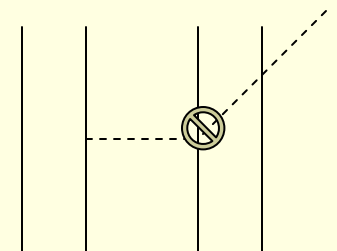
# $pp \rightarrow pp \pi^0$ results for s-wave $\pi$

- Same  ${}^3P_0 \rightarrow {}^1S_0$ s quantum numbers!
- Threshold s-wave cross section 5 times larger than one-body contribution.
- Neutral pion rescattering small on-shell.
- Large contributions from sigma and omega heavy meson exchange currents can explain data.
- Convergence of chiral pert theory poor because of large nucleon momenta even at threshold.

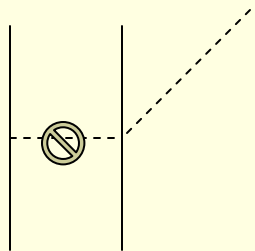
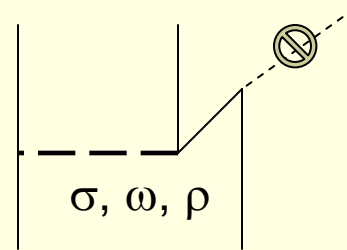
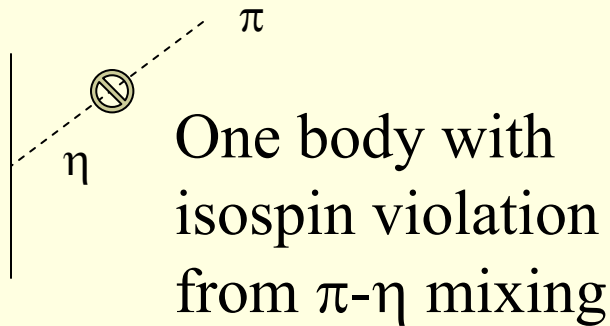


# Leading Order Contribution Near Threshold is Pion Rescattering Term.

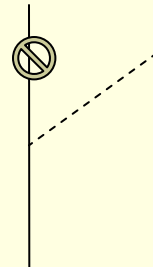
- Chiral symm relates isospin breaking mass and  $\pi^0$ -nucleon scattering terms
- $L_{qm} = \delta m / 2 \{ N^+ \tau_3 N - 2 N^+ \pi_3 \pi \cdot \tau N / F_\pi^2 \},$
- $L_{em} = \bar{\delta} m / 2 \{ N^+ \tau_3 N + 2 N^+ (\pi_3 \pi \cdot \tau - \pi^2 \tau_3) N / F_\pi^2 \},$
- $\Delta M = \delta m + \bar{\delta} m$
- Strong isospin dependence of these terms reduces their contribution to  $dd \rightarrow {}^4\text{He} \pi^0$
- Not true for CSB in  $np \rightarrow d \pi^0$  TRIUMF experiment



# N<sup>2</sup>LO Contributions

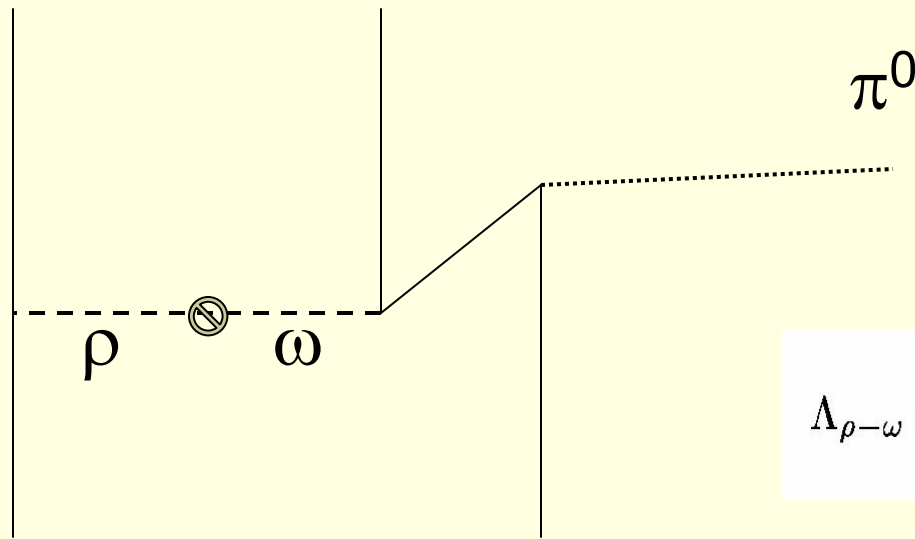


Pion rescattering after  $\pi$ - $\eta$  mixing



Isospin violation in  ${}^4\text{He}$  wave function, very sensitive to difference  $r_n - r_p$

# Also Rho-Omega Mixing



$$\Lambda_{\rho-\omega} = \frac{1}{(2m_\pi)^{1/2}} f\left(\frac{\langle \rho | H\omega \rangle}{m_\omega^2}\right) \frac{1}{M}$$

$$H_{\rho-\omega} = -\Lambda_{\rho-\omega} \left(\frac{g_\rho g_\omega}{4\pi M}\right) \frac{1}{2} \sum_{i \neq j} \left\{ (1 + \tau_i^3 \tau_j^3) \sigma_j \cdot (f_{ij}^\rho p_i + p_i f_{ij}^\rho) + i[1 + \tau_i^3 \tau_j^3 (1 + C_\rho)] \sigma_i \wedge \sigma_j \cdot q_j f_{ij}^\rho \right\}$$



# Plane wave calculation to characterize operators

## Anti-symmetrized spin-isospin wave functions

The spin-isospin parts of the wave functions are

$$|\alpha\rangle = \frac{1}{\sqrt{2}} \left( [(12)_1(34)_1]_0 [(12)_0(34)_0]_0 \right. \\ \left. - [(12)_0(34)_0]_0 [(12)_1(34)_1]_0 \right)$$

$$|d_1\rangle = (12)_1(12)_0$$

$$|d_2\rangle = (34)_1(34)_0$$

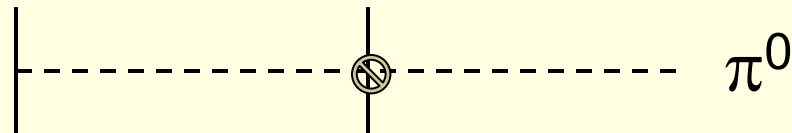
where the first brackets is spin and the second isospin.

# $\pi$ Rescattering is Isospin Disfavored

- Operator has two isospins ( $\tau$ ) but only one spin operator ( $\sigma$ ).

$$\mathcal{O}_{i,j} = -\frac{f}{4\pi\mu F_\pi^2 \sqrt{2\mu}} [\delta m_N (\boldsymbol{\tau}^{(i)} \cdot \boldsymbol{\tau}^{(j)} + \tau_3^{(i)} \tau_3^{(j)}) - \bar{\delta} m_N (\boldsymbol{\tau}^{(i)} \cdot \boldsymbol{\tau}^{(j)} - \tau_3^{(i)} \tau_3^{(j)})] \\ (\boldsymbol{\sigma}^{(i)} - \boldsymbol{\sigma}^{(j)}) \cdot (-i\hat{\mathbf{r}}) \frac{d}{dr} \left( \frac{e^{-\mu r}}{r} \right),$$

- The two  $\tau$ s convert both deuterons to isospin 1 but  $\sigma$  can only convert one of the two deuterons spins to spin 0.
- The s-wave alpha particle is a mixture of spin 0, isospin 1 pairs or spin one, isospin zero pairs.
- Operator gives zero for plane wave initial state. Distortions and or d-waves in  ${}^4\text{He}$  or deuterons can change this.



# Other Contributions add coherently

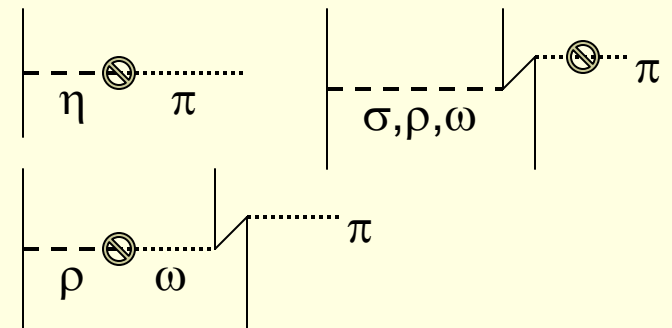
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- 1-body and sigma are coherent
- Omega involves  $-\sigma_3 \cdot p_1$  compared to  $\sigma_1 \cdot p_1$  for sigma exchange. For  ${}^3P_0 \rightarrow {}^1S_0$ ,  $\sigma_3 \cdot p_1 = -\sigma_1 \cdot p_1$  so omega and sigma add coherently.
- One body +  $\sigma + \omega$  HMEC are coherent.
- Should evaluate rho-omega mixing and other HME in  $np \rightarrow d\pi^0$ .

# Preliminary results with gaussian He and d and plane wave dd initial state

- Pure s-wave  $\pi$  at lower  $E=228.5$  MeV
- One body  $\pi$ - $\eta$  mixing alone  $\sigma=4.8$  pb [0.8 pb].
- Total cross section  $\sigma_{1+\sigma+\omega+\rho+\rho-\omega}=37.7$  pb [8.7 pb]
- Experiment=  $12.7 \pm 2.2$  pb
- Based on  $\langle \pi H \eta \rangle = -4200$  MeV<sup>2</sup>,  $g_\eta^2/4\pi = 3.68$  [0.6] and  $\langle \rho H \omega \rangle = -4500$  MeV<sup>2</sup>.

Heavy Meson Exc. Amplitude	Comp. to 1-Body
$\sigma$	0.62
$\omega$	0.66
$\rho$	0.17
$\rho$ - $\omega$	0.36



# Coherent Pion Bremsstrahlung

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- Can arrange spin one deuterons so axial charges of all 4 nucleons contribute coherently.
- We find heavy meson exchange currents increase cross section because all pairs contribute with same sign and add coherently to one-body.
- Plane wave calculations with heavy meson exchange and isospin violation from reasonable  $\pi$ - $\eta$  mixing parameters can explain large observed cross section.
- Important to include dd distortions and full  ${}^4\text{He}$  wave function and to calculate Weinberg terms with correlations! Calculations underway.

# CSB in $np \rightarrow d\pi^0$ (A. K. Opper)

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- Forward, backward asymmetry violates CS

$$A_{fb}(\theta) = (\sigma(\theta) - \sigma(\pi - \theta)) / (\sigma(\theta) + \sigma(\pi - \theta))$$

- TRIUMF exp at 279.5 MeV detects both forward and backward d in same spectrometer setting. Result angle integrated:

$$A_{fb} = +17.2 \pm 8(\text{stat}) \pm 5.5(\text{sys}) \times 10^{-4}$$

based on careful simulation of spectrometer acceptance.

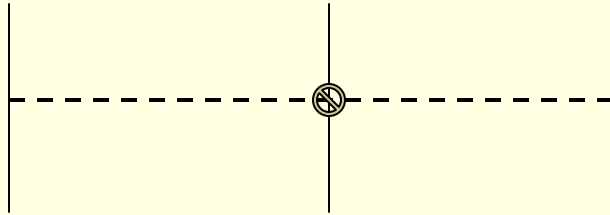
# $np \rightarrow d\pi^0$ Theory (Niskanen)

- $A_{fb}$  depends on interference of s+p waves and CSB and CS amplitudes. This increases uncertainty. No isospin cancellations for  $\delta m$ ,  $\bar{\delta m}$  terms:  
$$A_{fb} = -28 \left\{ g_{\eta NN} / (4\pi(3.68))^{1/2} \langle \pi H\eta \rangle / (-5900 \text{ MeV}^2) - 0.87/\text{MeV}(\delta m - \bar{\delta m}/2) \right\} \times 10^{-4}.$$

This is – for  $\pi\eta$  mixing, + for  $\delta m$ .

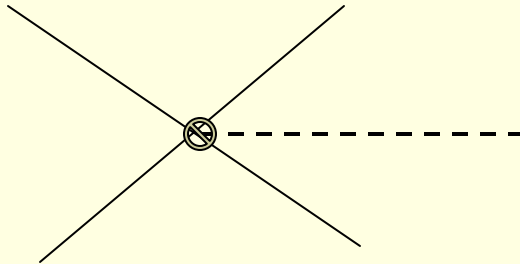
- Positive exp result may be evidence for  $\delta m$  terms.
- Different isospin dependence makes  $np \rightarrow d\pi^0$  and  $dd \rightarrow {}^4\text{He}\pi^0$  sensitive to different combinations of CSB operators.
- Important to improve calculation of  $A_{fb}$ , including HME for example, and assess theoretical uncert.

# CSB Operators in Chiral Pert Theory



Contributes to  $np \rightarrow d\pi^0$ ,  
isospin suppressed for  
 $dd \rightarrow {}^4\text{He } \pi^0$

Leading order:  $(\sigma_1 - \sigma_2) \cdot q \tau_1 \tau_2$



Important for  
 $dd \rightarrow {}^4\text{He } \pi^0$ ,  
should be  
calculated for  
 $np \rightarrow d\pi^0$

Contact term: down by  $m_\pi/M \sim 1/7$

$\sigma_1 \cdot p$  for  $\sigma, \omega$  HMEC or  $\sigma_1 \times \sigma_2 \cdot q \tau_1 \tau_2$  for  $\rho - \omega$